

UNRAVELLING THE RESONANT INSTABILITIES OF A STRATIFIED GRAVITY WAVE

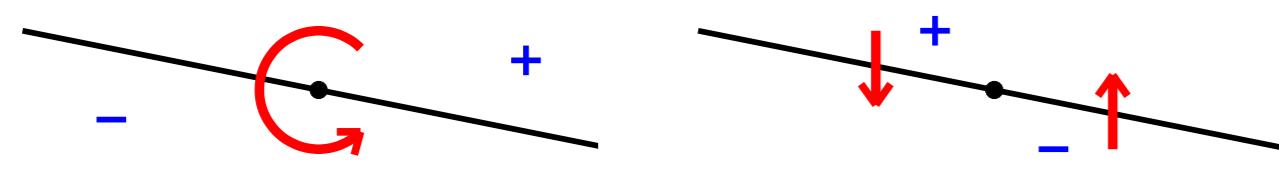
YUANXUN BILL BAO & DAVID J. MURAKI

DEPARTMENT OF MATHEMATICS, SIMON FRASER UNIVERSITY, BURNABY, BC, CANADA

WAVES IN A 2D STRATIFIED FLUID

Boussinesq Fluid with Constant Stratification

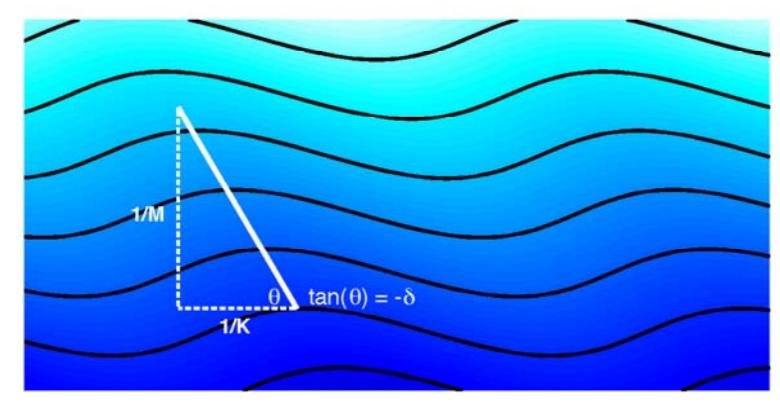
$$\nabla \cdot \vec{u} = 0 ; \quad \frac{D\eta}{Dt} = -b_x ; \quad \frac{Db}{Dt} = -N^2 w$$



- Streamfunction, $\psi(x, z, t)$ & velocity, $\vec{u} = (u, w) = (\psi_z, -\psi_x)$
- Buoyancy, $b(x, z, t)$ & vorticity, $\eta(x, z, t) = \psi_{zz} + \psi_{xx}$
- Brunt-Väisälä frequency, N & stable stratification

Exact Nonlinear Solution — Finite-Amplitude Gravity Wave

$$\begin{pmatrix} \psi \\ b \end{pmatrix} = \begin{pmatrix} -\Omega/KM \\ N^2/M \end{pmatrix} 2\epsilon \sin(Kx + Mz - \Omega t)$$

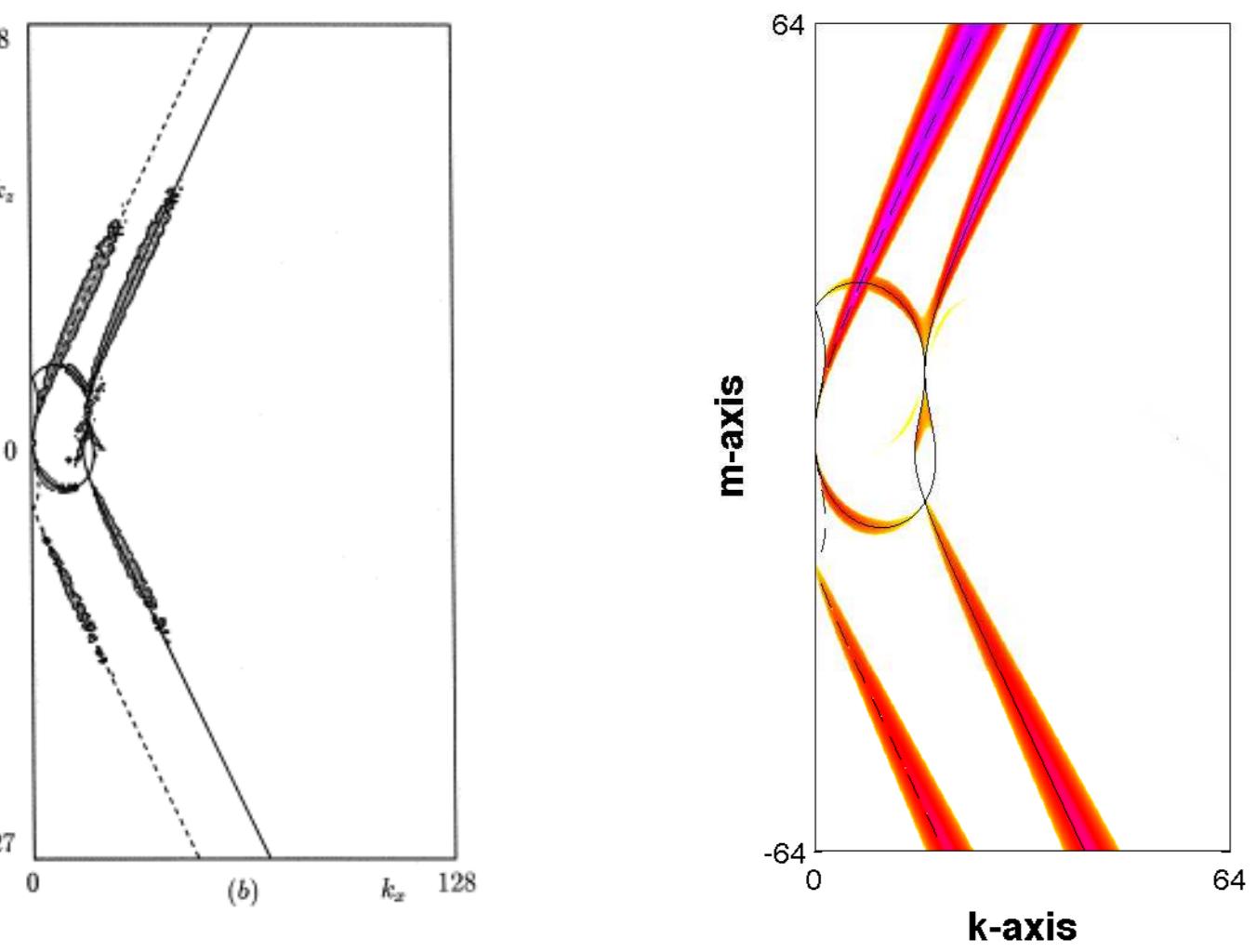


- Primary wavenumbers: (K, M)
- Propagation angle: $\delta = -\frac{K}{M}$

Dispersion relation: $\Omega^2(K, M) = \frac{N^2 K^2}{K^2 + M^2}$

UNSTABLE SPECTRUM: DNS & FLOQUET

- Linear instability with wavenumbers (k, m)
- Growing modes from Direct Numerical Simulations (Lin, 2000) — (left)
- Motivation: unstable eigenvalues using Floquet theory — (right)



References

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LINEAR STABILITY ANALYSIS

- Linear stability analysis on dimensionless gravity wave + disturbances

$$\begin{pmatrix} \psi \\ b \end{pmatrix} = \begin{pmatrix} -\Omega \\ 1 \end{pmatrix} 2\epsilon \sin(x + z - \Omega t) + \begin{pmatrix} \tilde{\psi}(x, z, t) \\ \tilde{b}(x, z, t) \end{pmatrix}$$

Dimensionless frequency: $\Omega^2 = \frac{1}{1 + \delta^2}$

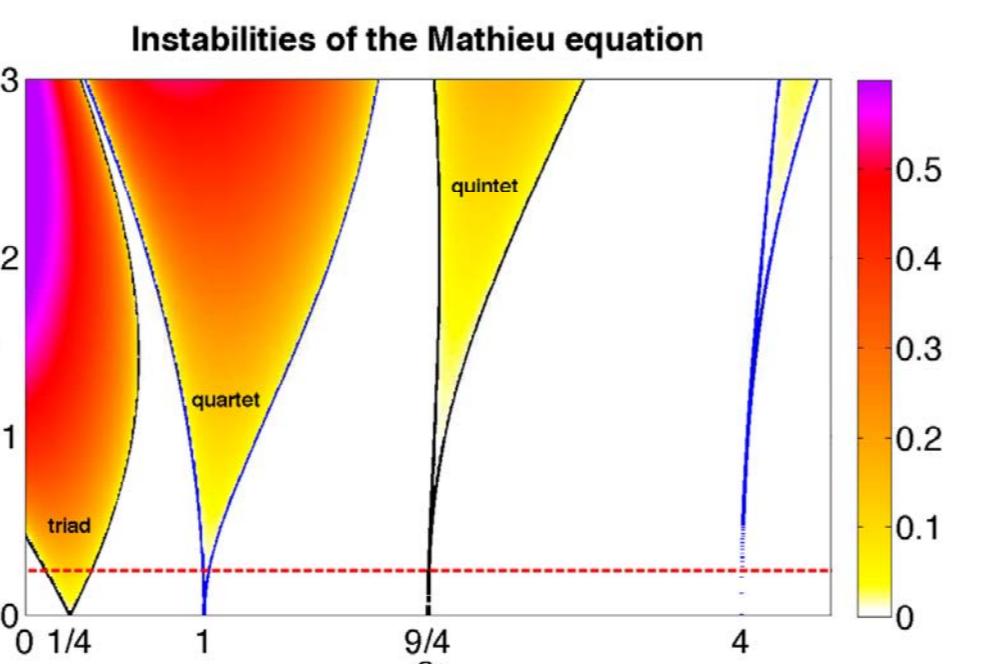
- Linear PDE system with non-constant, but periodic coefficients

$$\begin{aligned} \tilde{\eta}_t + \tilde{b}_x - 2\epsilon J(\Omega\tilde{\eta} + \tilde{\psi}/\Omega, \sin(x + z - \Omega t)) &= 0 \\ \tilde{b}_t - \tilde{\psi}_x - 2\epsilon J(-\Omega\tilde{b} + \tilde{\psi}, \sin(x + z - \Omega t)) &= 0 \\ \tilde{\psi}_{zz} + \delta^2 \tilde{\psi}_{xx} &= \tilde{\eta} \end{aligned}$$

- Linear advection in Jacobian:

$$J(f, g) = \begin{vmatrix} f_x & g_x \\ f_z & g_z \end{vmatrix} = f_x g_z - g_x f_z$$

FLOQUET THEORY

Instabilities of the Mathieu Equation: $\ddot{u} + (\alpha + \epsilon \sin t) u = 0$ 

- Floquet representation with Fourier series:

$$u(t) = e^{-i\omega t} \left\{ \sum_{-\infty}^{+\infty} v_n e^{-int} \right\} = \text{exponential part} \times \left\{ \text{co-periodic part} \right\}$$

- Floquet exponent, $\omega(\alpha; \epsilon)$ & $\text{Im } \omega > 0 \rightarrow$ instability

Floquet/Fourier Analysis for PDEs

- Floquet representation with disturbance wavenumbers, (k, m)

$$\begin{pmatrix} \tilde{\psi} \\ \tilde{b} \end{pmatrix} = e^{i(kx + mz - \omega t)} \left\{ \sum_{-\infty}^{+\infty} \tilde{v}_n e^{in(x+z-\Omega t)} \right\}$$

- Floquet exponent, $\omega(k, m; \epsilon)$ & $\text{Im } \omega > 0 \rightarrow$ instability

Dispersion relation for disturbances: $\omega^\pm(k, m; 0) = \pm \frac{|k|}{\sqrt{\delta^2 k^2 + m^2}}$

- A generalized eigenvalue problem with Hill's infinite matrix

$$\begin{bmatrix} \dots & \dots & & \\ \dots & S_0 & \epsilon M_1 & \\ \epsilon M_0 & S_1 & \dots & \dots \end{bmatrix} - \omega \begin{bmatrix} \dots & & & \\ & \Lambda_0 & & \\ & & \Lambda_1 & \\ & & & \dots \end{bmatrix}$$

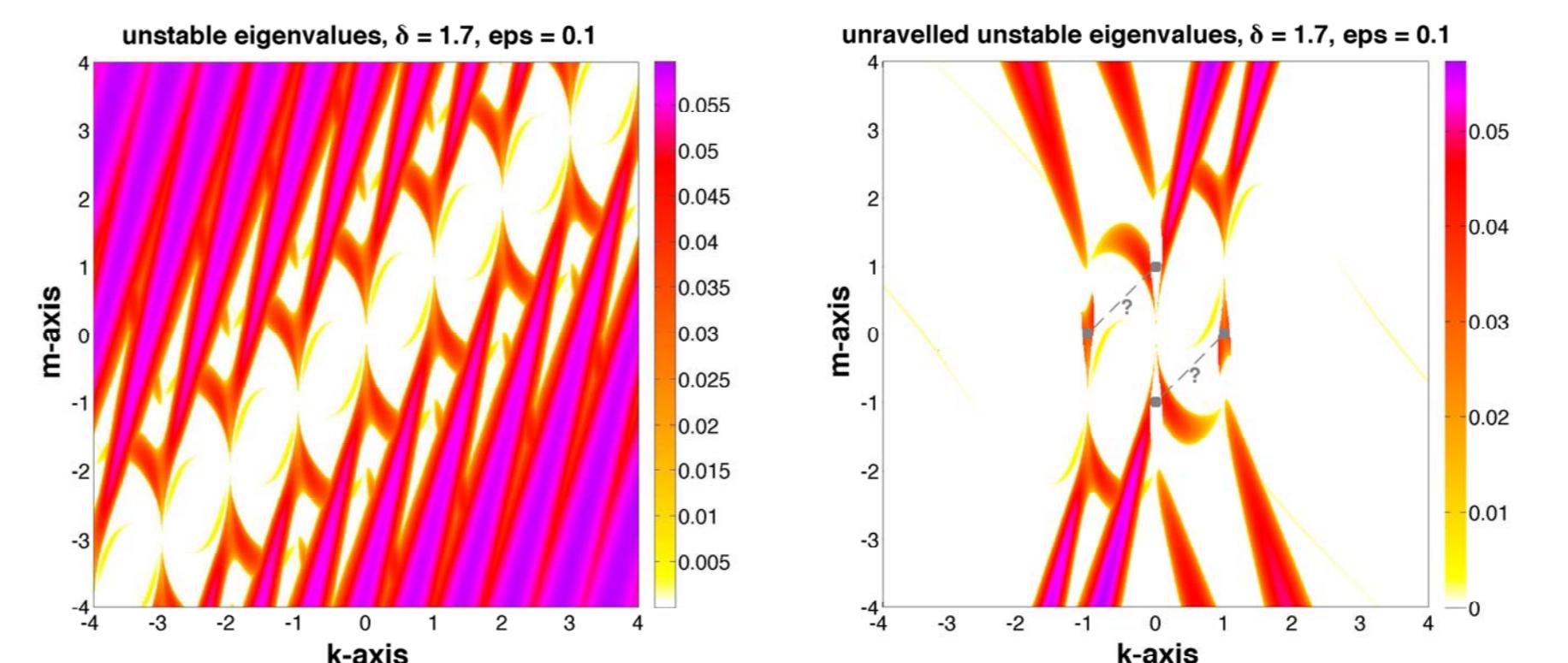
- $2 \times 2 \mathbb{R}$ blocks: $M_n(k, m)$; $S_n(k, m)$, symmetric; $\Lambda_n(k, m)$, diagonal
- Truncate $-N \leq n \leq N$ & compute $4N+2$ eigenvalues $\{\omega(k, m; \epsilon)\}$

A TANGLE OF UNSTABLE EIGENVALUES

- Unstable Floquet eigenvalues selected by maximum growth rate — (left)

- Periodicity from index shifts \rightarrow multiple counting of $\text{Im } \omega$'s

$$\begin{pmatrix} \tilde{\psi} \\ \tilde{b} \end{pmatrix} = e^{i((k+q)x + (m+q)z - (\omega + \Omega q)t)} \left\{ \sum_{-\infty}^{+\infty} \tilde{v}_n e^{in(x+z-\Omega t)} \right\}$$



- QUESTION 1: is there an association of $\omega(k, m; \epsilon)$ with instabilities corresponding to physical wave resonances — as in Lin (2000)?

- Q1: Yes! An unravelling of unstable Floquet eigenvalues — (right)

- QUESTION 2: are there two values of $\omega(k, m; \epsilon)$ as in the dispersion relation, $\omega^\pm(k, m; 0)$; or $4N+2$ as from the truncated Hill's matrix?

- QUESTION 3: is the index-periodicity of the Floquet/Fourier method a numerical artifact, or a natural feature of Floquet theory?

PERTURBATION ANALYSIS

- Two branches from ϵ -perturbation theory: $\omega^\pm(k, m; \epsilon) \sim \omega^\pm(k, m; 0)$

- Complex ω 's arise from ϵ -perturbation of multiple Hill's eigenvalues

Triad ($n = \pm 1$) Resonant Trace Curves

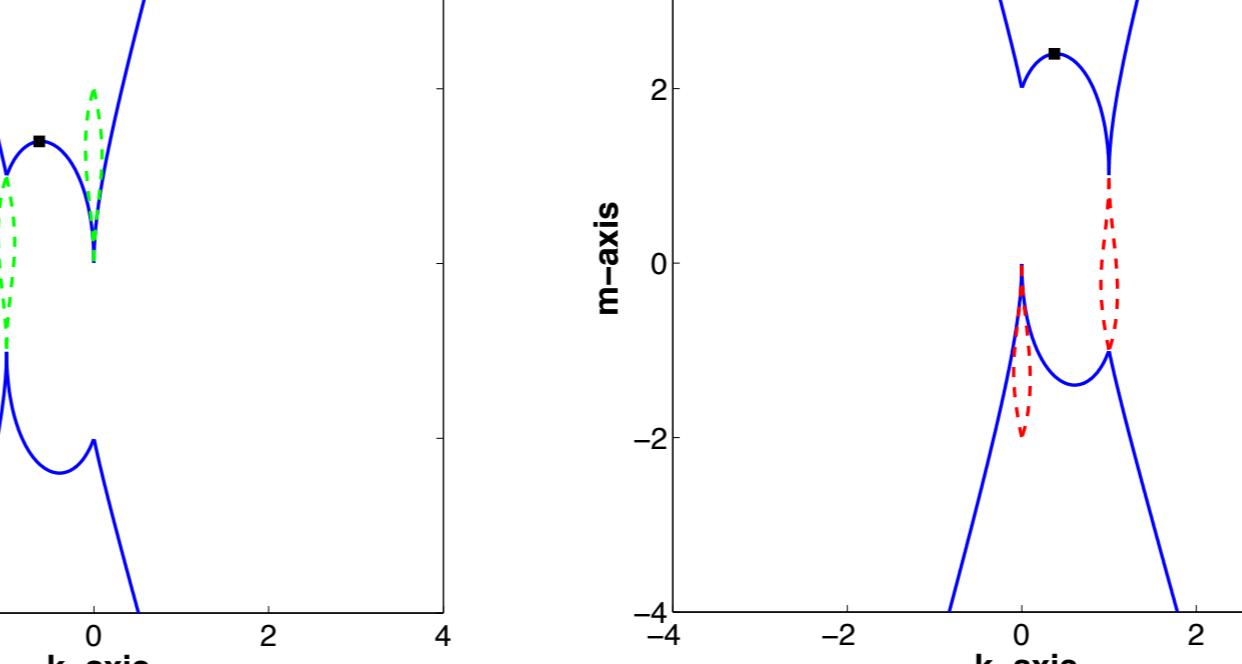
- (k, m) — trace curve for triad resonances: (ω^+, left) & (ω^-, right)

$$\omega^\pm(k, m; 0) = \omega^\pm(k + n, m + n; 0) - n\Omega$$

- Triad trace curves are also where $\omega^\pm(k, m; 0)$ are double eigenvalues

$\omega^+(k, m; 0)$ resonant traces

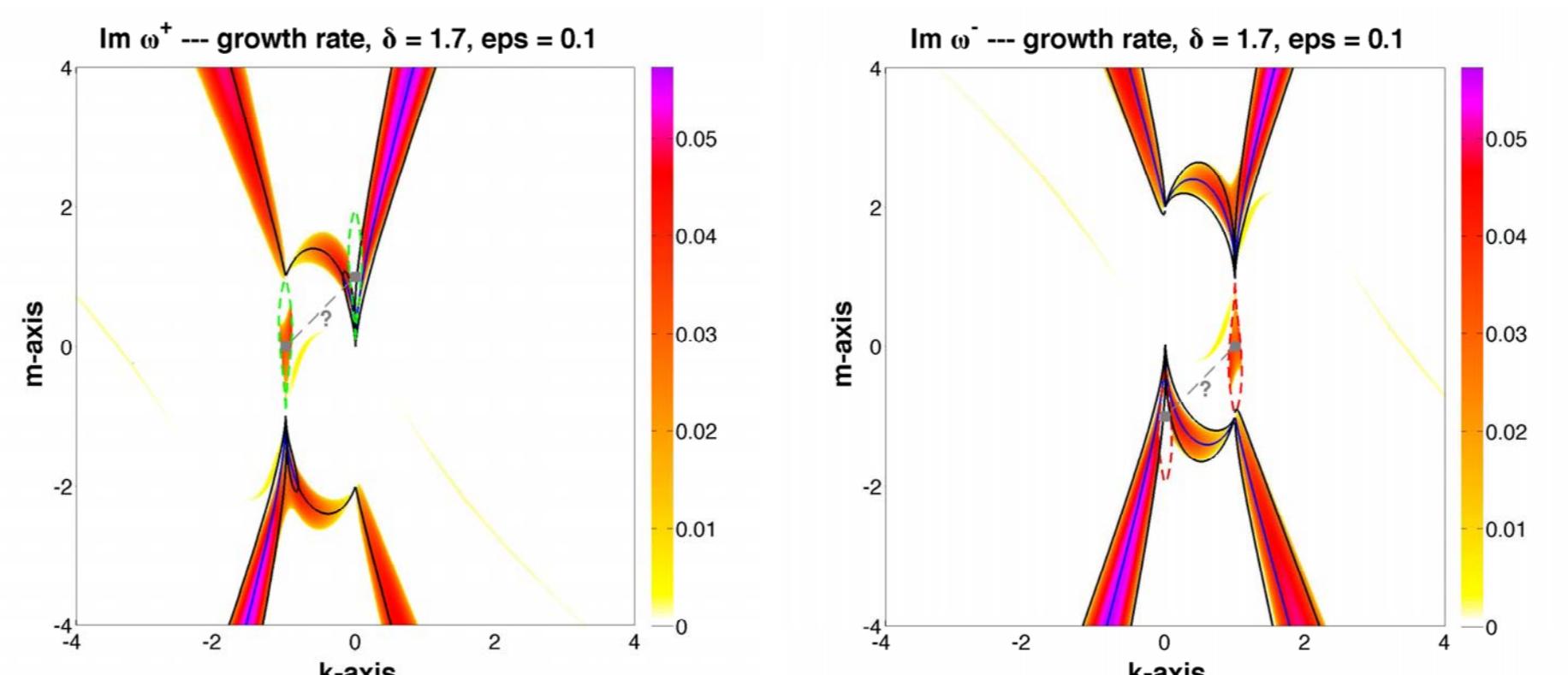
$\omega^-(k, m; 0)$ resonant traces



- Triad blue traces are unstable by perturbation theory

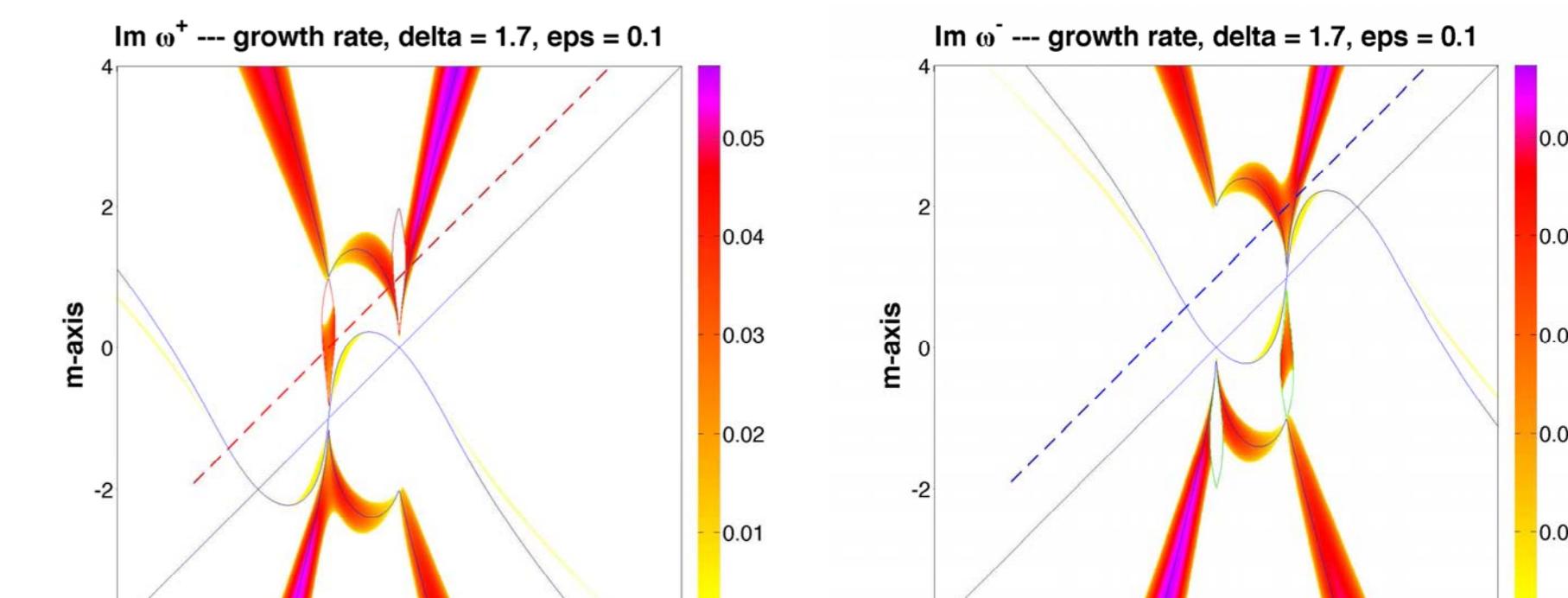
- Unravelled $\text{Im } \omega^\pm$ by ϵ -continuity: computational perturbation theory

- Stability boundaries (black) from analytical perturbation theory



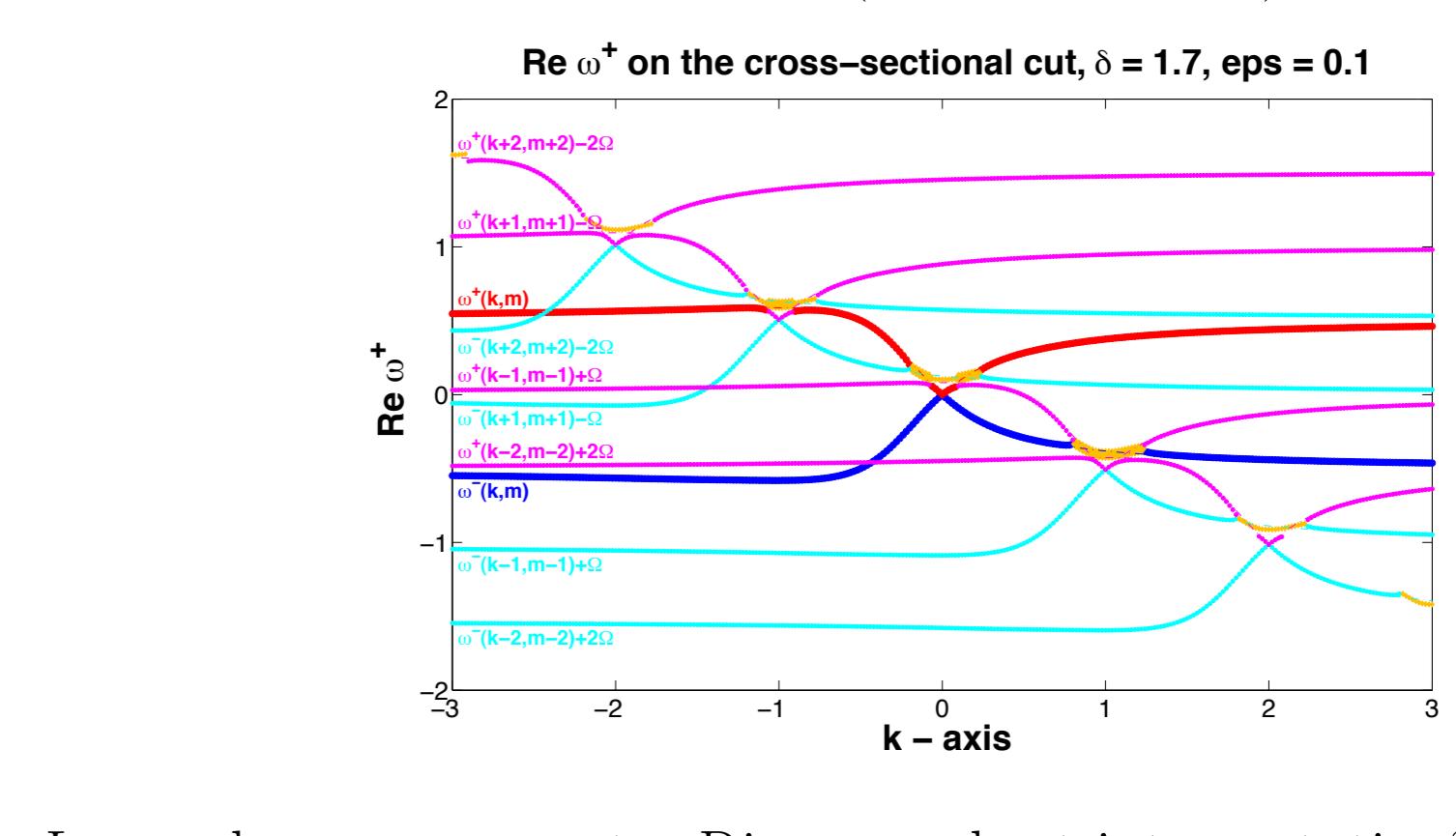
A RIEMANN SHEET PERSPECTIVE ?

- Plot 10 curves of $\text{Re } \omega^\pm$ along the cross-sectional cuts



- Unravelled values of $\text{Re } \omega^+(k, m)$ & $\text{Re } \omega^-(k, m)$

- Other curves correspond to $\omega^\pm(k + n, m + n; \epsilon) - n\Omega$

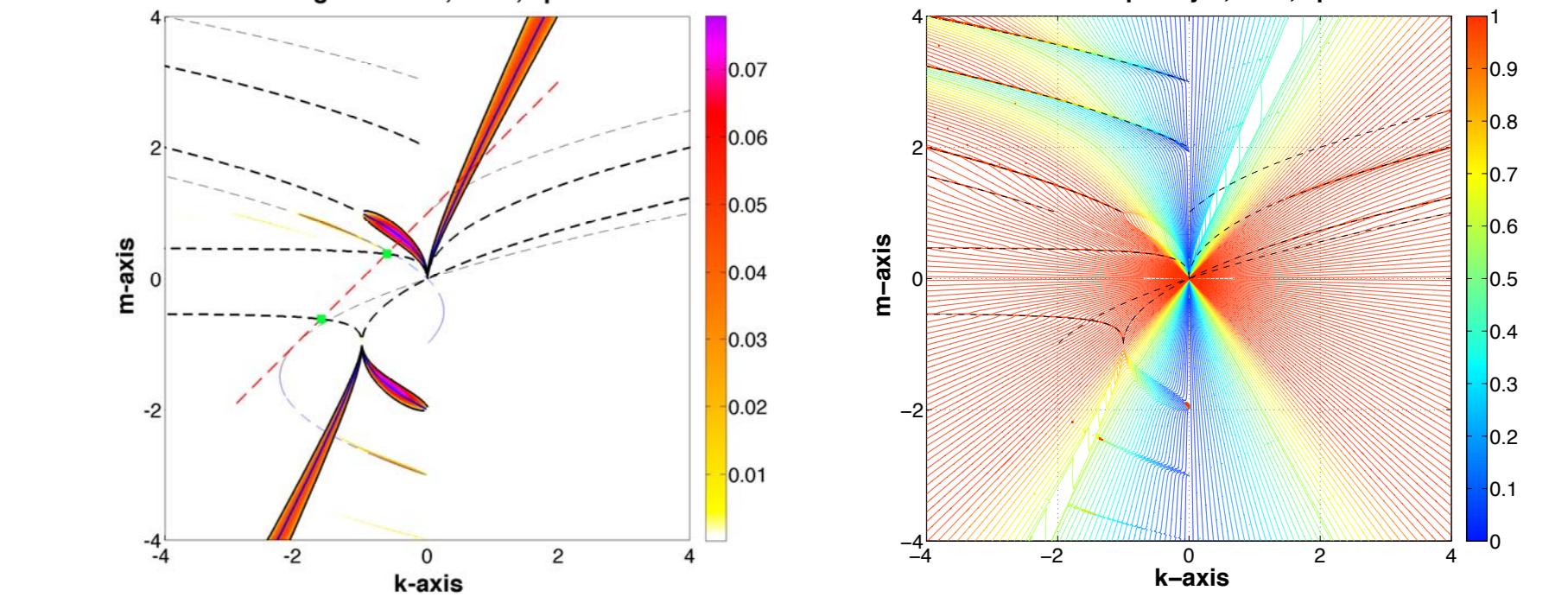


- Layered curves suggest a Riemann sheet interpretation?

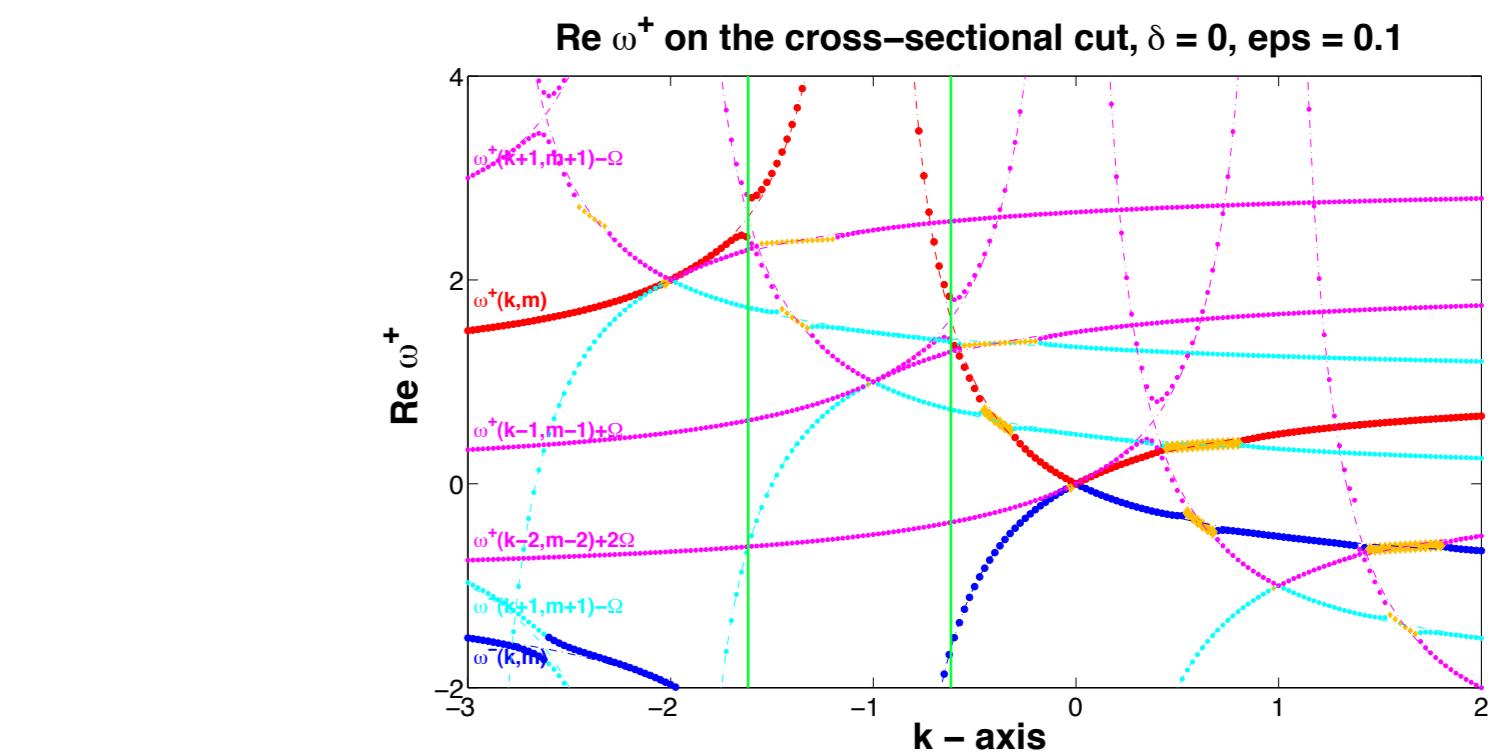
HYDROSTATIC LIMIT ($\delta = 0$)

- Unravelled $\text{Im } \omega^+$ with triad (thick) & quartet (thin) traces — (left)

- Frequency $\text{Re } \omega^+$ displays discontinuities across stable traces — (right)



- Discontinuities in $\text{Re } \omega^+$ suggest branch cuts in the Riemann sheets



- Connection: do Floquet/Fourier results reflect the Riemann sheets idea of McKean, et al?

PRELIMINARY CONJECTURES

- Q2: two values of $\omega^\pm(k, m; \epsilon)$ are physically important Riemann sheets

- Q2/Q3: $4N+2$ $\omega(k, m; \epsilon)$ reflect truncation of Riemann surface

- ?: complications at high-multiplicity eigenvalues \rightarrow crossing traces